

The background features a repeating pattern of interconnected, faceted shapes in shades of blue and red, resembling a molecular or crystal lattice structure. A dark blue horizontal band is overlaid across the middle of the image, containing the title text in white.

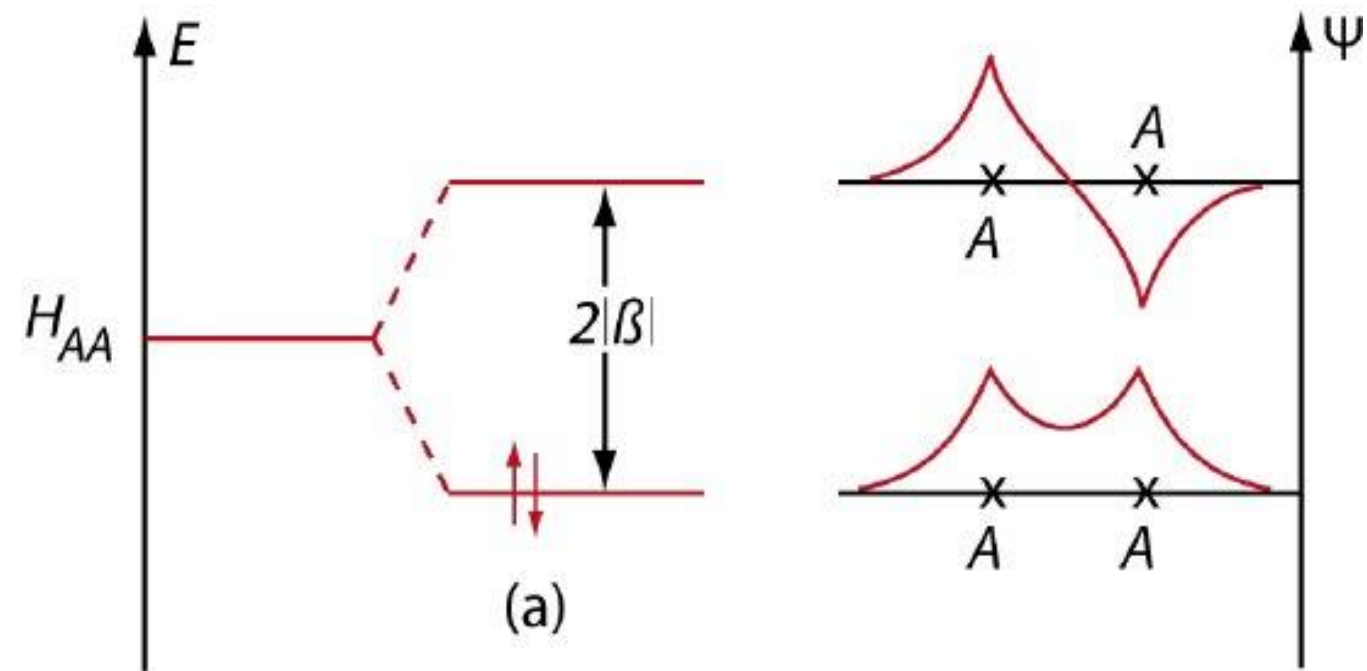
# Modifying the bonding character of coupled states of thin-plate elastic resonators via prestress modulation

Pragalv Karki<sup>1</sup> and Jayson Paulose<sup>2</sup>

1. Department of Radiology, Mayo Clinic, Rochester, MN, USA
2. Department of Physics, University of Oregon, OR, USA

# Bonding character

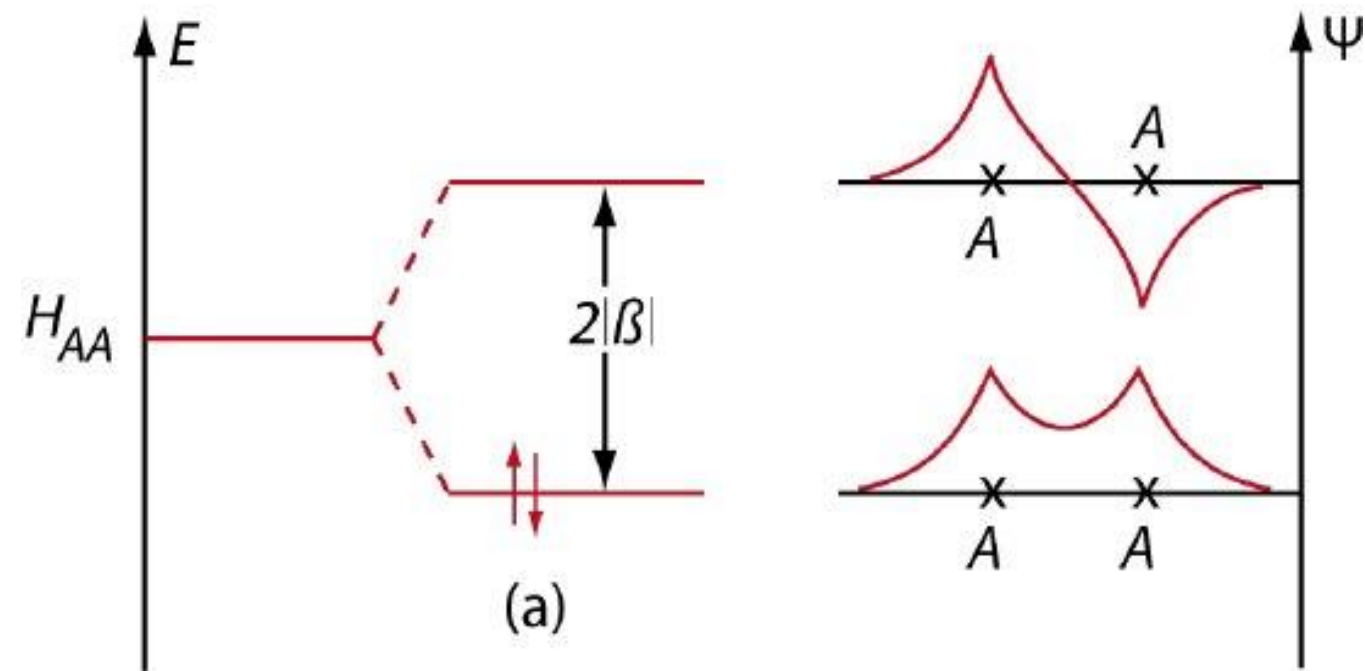
Diatomic molecule with expected bonding character



D. G. Pettifor, arXiv:1112.4638

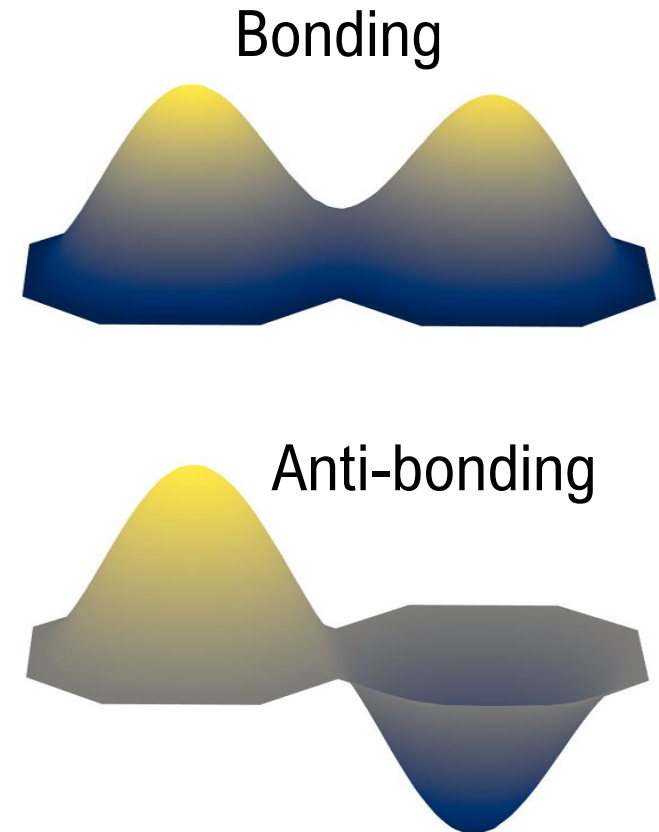
# Bonding character

Diatomic molecule with expected bonding character



D. G. Pettifor, arXiv:1112.4638

Reversing the bonding character in thin elastic plates

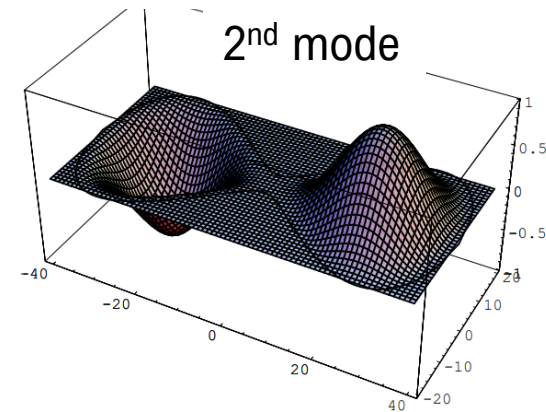
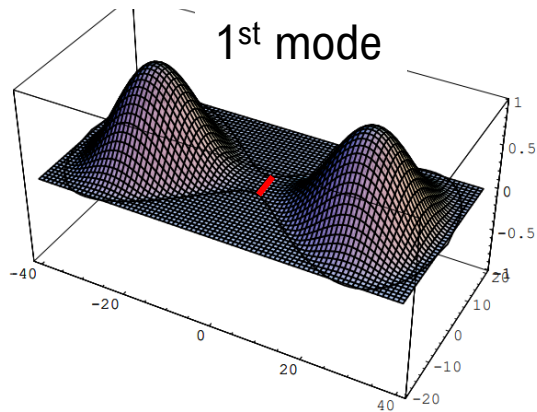
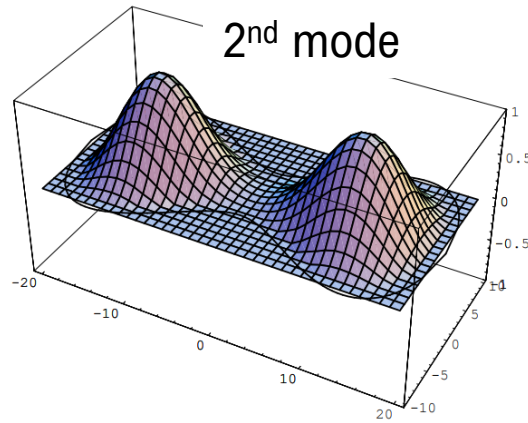
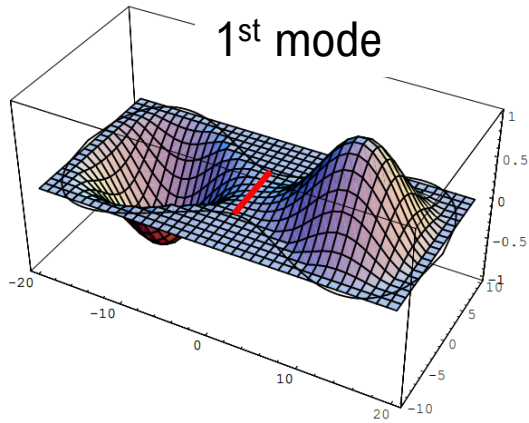


# Fourth order plate to thin elastic plate

## Fourth order plate

$$\nabla^4 u = 0 \quad \text{on domain}$$

$$u = \nabla u = 0 \quad \text{clamped on boundary}$$

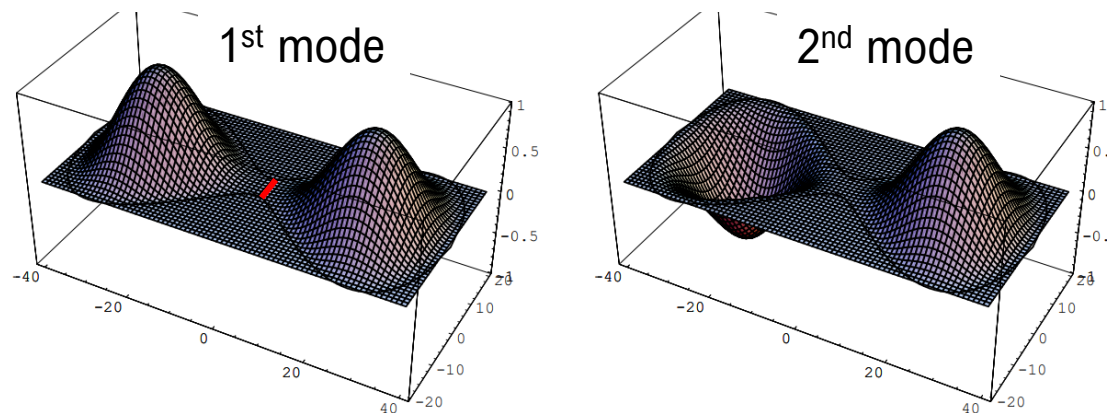
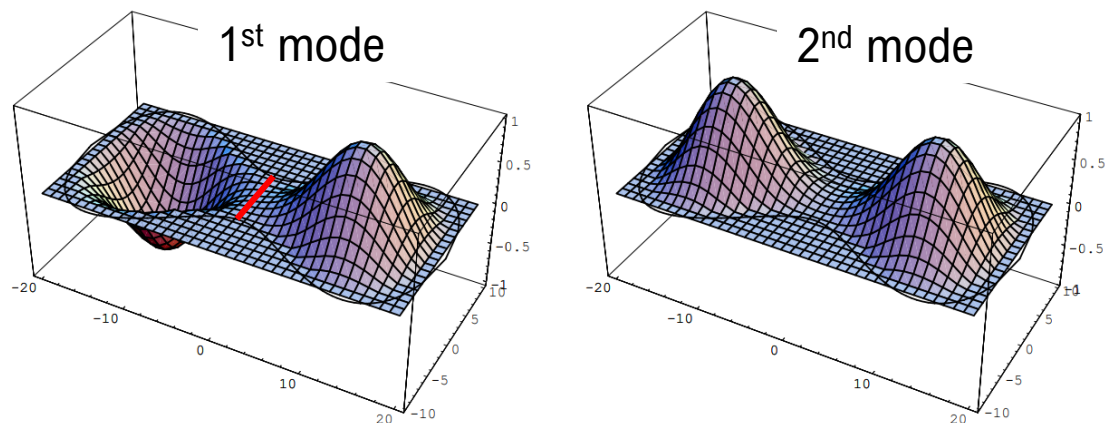


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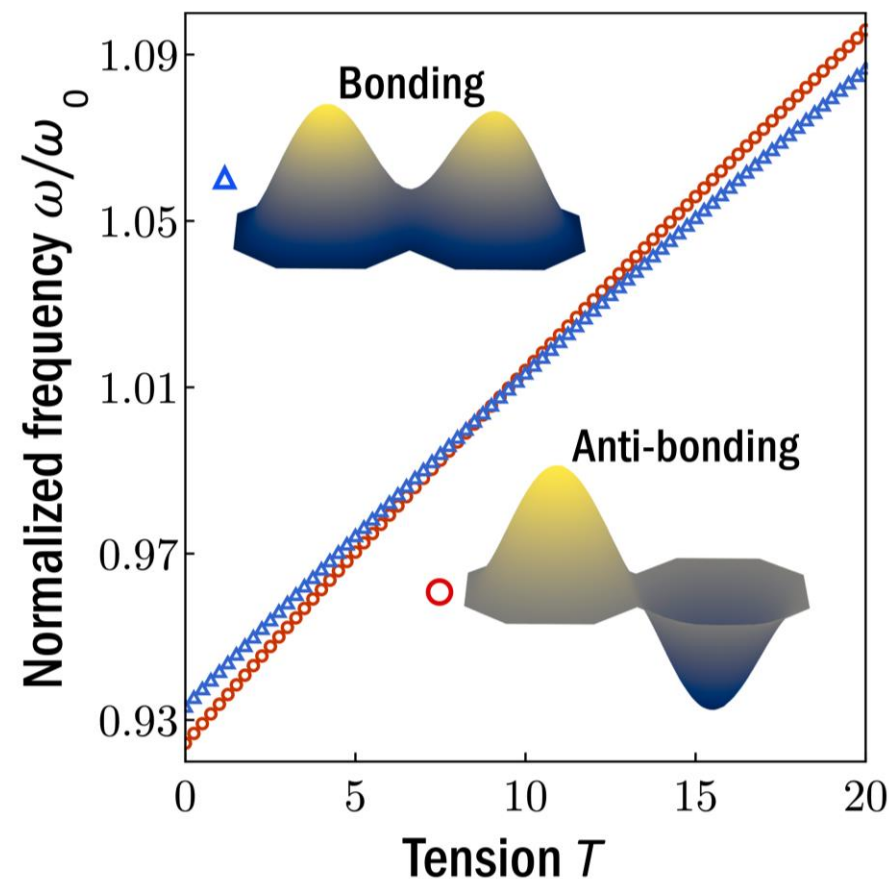
$$u = \nabla u = 0 \quad \text{clamped on boundary}$$



## Thin elastic plate (Föppl–von Kármán)

$$\nabla^4 u - T\nabla^2 u = 0 \quad \text{on domain}$$

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# Thin elastic plate (Föppl–von Kármán)

Full dynamical equation

$$\rho \frac{\partial^2 u}{\partial t^2} + c \frac{\partial u}{\partial t} + D \nabla^4 u - T' \nabla^2 u = 0$$

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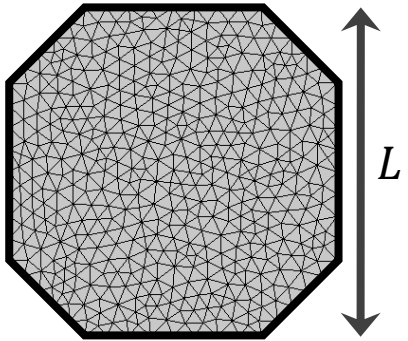
$$\rho \frac{\partial^2 u}{\partial t^2} + c \frac{\partial u}{\partial t} + D \nabla^4 u - T' \nabla^2 u = 0$$

Rescale

$$\bar{x} = x/L, \bar{y} = y/L, \bar{t} = t \sqrt{D/(\rho L^4)}$$

Non-dimensional form

$$\frac{\partial^2 u}{\partial \bar{t}^2} + \zeta \frac{\partial u}{\partial \bar{t}} + \bar{\nabla}^4 u - T \bar{\nabla}^2 u = 0$$



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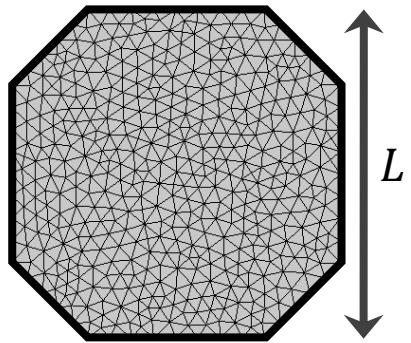
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Drop the bars

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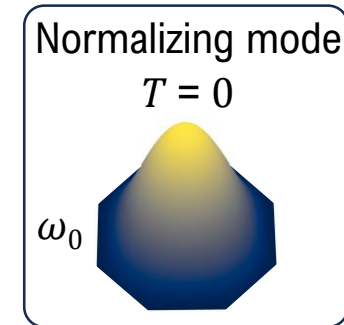
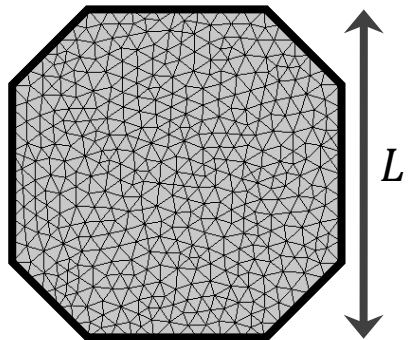
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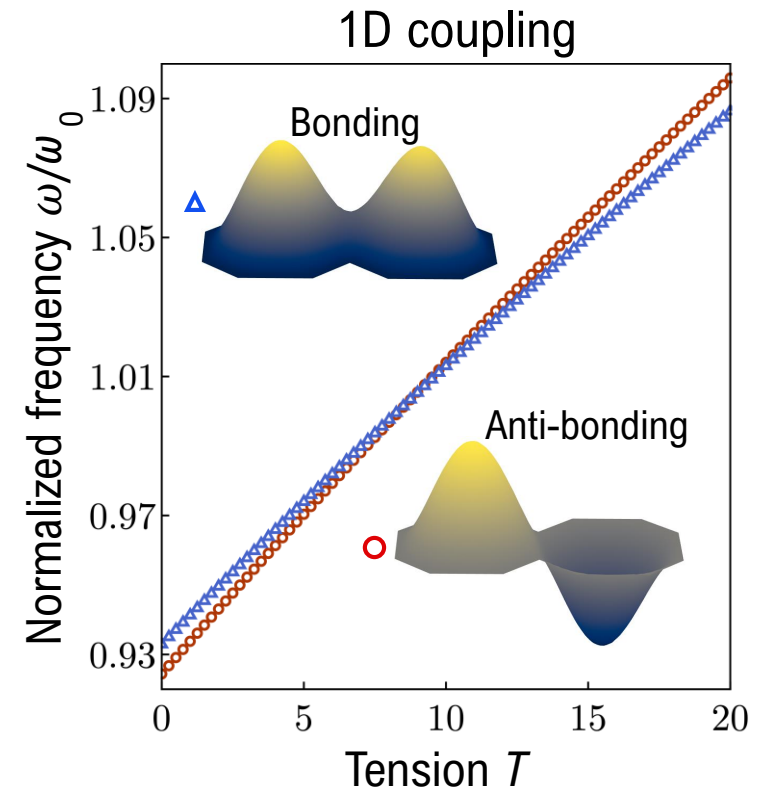
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Eigenvalue problem

$$\nabla^4 u - T \nabla^2 u = \omega^2 u$$



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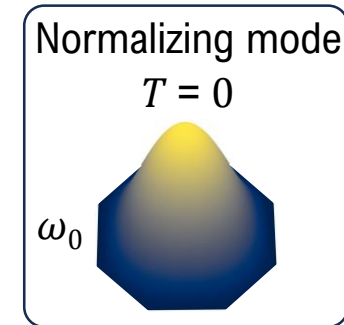
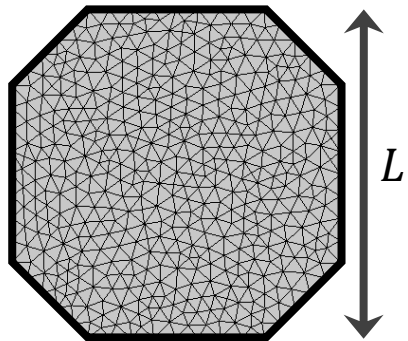
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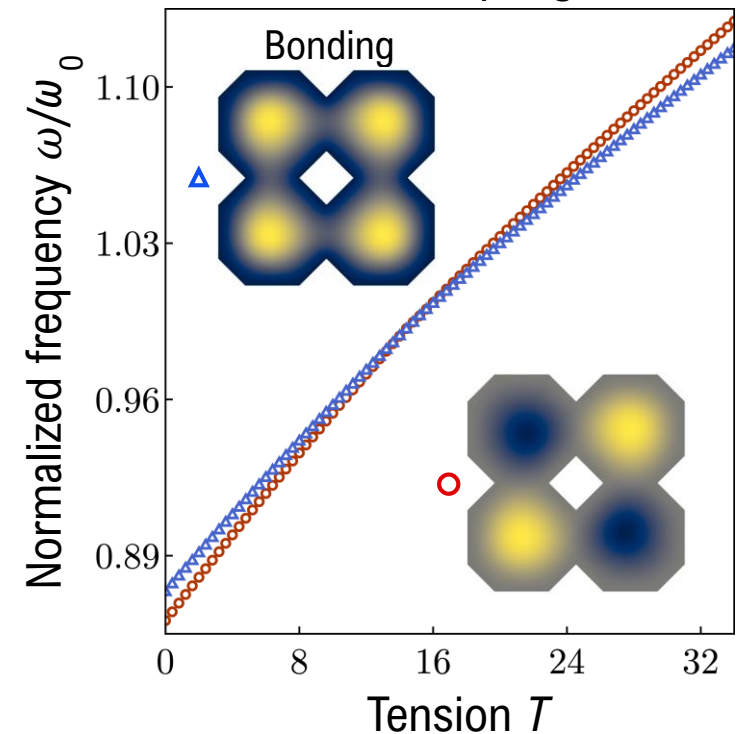
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2D coupling



# COMSOL implementation

## Eigenvalue problem

$$\nabla^4 u - T\nabla^2 u = \omega^2 u$$

## Eigenvalue Study in General Form PDE

$$\lambda^2 e_a \mathbf{u} - \lambda d_a \mathbf{u} + \nabla \cdot \Gamma = f$$

$$\mathbf{u} = [u1, u2, u3, u4, u5]^T$$

$$\nabla = \left[ \frac{\partial}{\partial x}, \frac{\partial}{\partial y} \right]$$

$$\Gamma_1 = \begin{bmatrix} u4_x + 2u5_x - Tu1_x \\ u1 \\ 0 \\ u2 \\ 0 \end{bmatrix} \quad \Gamma_2 = \begin{bmatrix} u5_y - Tu1_y \\ 0 \\ u1 \\ 0 \\ u3 \end{bmatrix} \quad f = \begin{bmatrix} 0 \\ u2 \\ u3 \\ u4 \\ u5 \end{bmatrix}$$

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## We get 5 sets of equations

$$\text{I) } (u4_{xx} + 2u5_{xx} + u5_{yy}) - T(u1_{xx} + u1_{yy}) = \lambda u1$$

$$\text{II) } u2 = u1_x$$

$$\text{III) } u3 = u1_y$$

$$\text{IV) } u4 = u2_x$$

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## Which simplifies to

$$\left( \frac{\partial^4 u1}{\partial x^4} + 2 \frac{\partial^4 u1}{\partial x^2 \partial y^2} + \frac{\partial^4 u1}{\partial y^4} \right) - T \left( \frac{\partial^2 u1}{\partial x^2} + \frac{\partial^2 u1}{\partial y^2} \right) = \lambda u1,$$

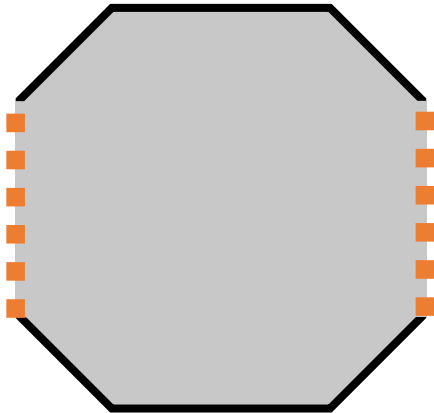
$$\nabla^4 u1 - T\nabla^2 u1 = \lambda u1$$

# Band structure calculation

Bloch periodicity:

$$u_{k_x}(x, y) = e^{ik_x x} \phi_{k_x}(x, y)$$

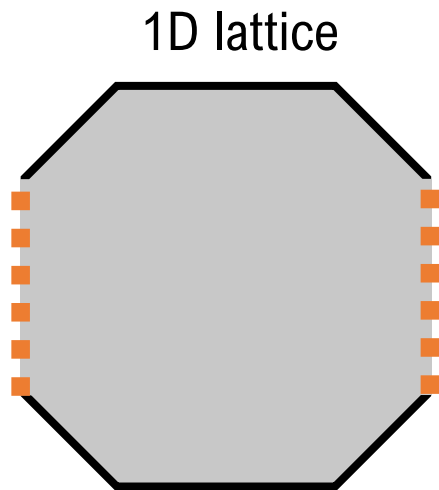
1D lattice



# Band structure calculation

Bloch periodicity:

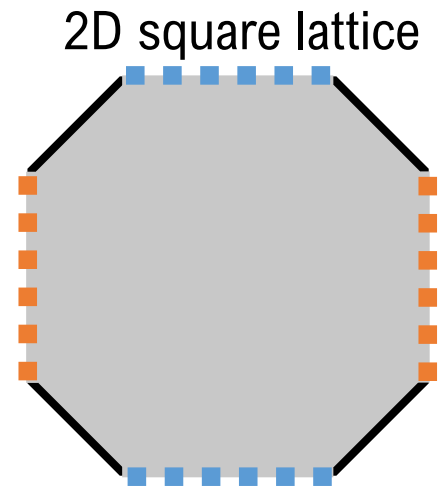
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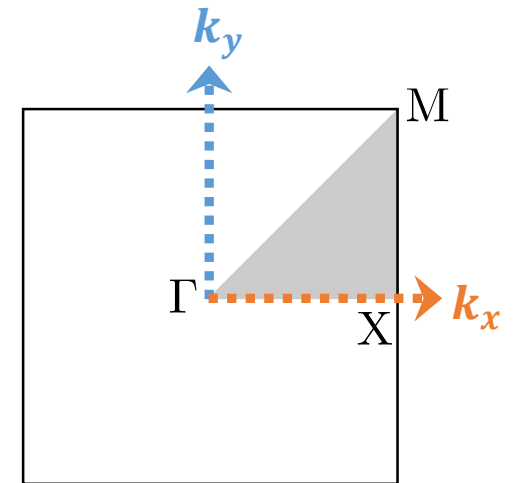
Bloch periodicity:

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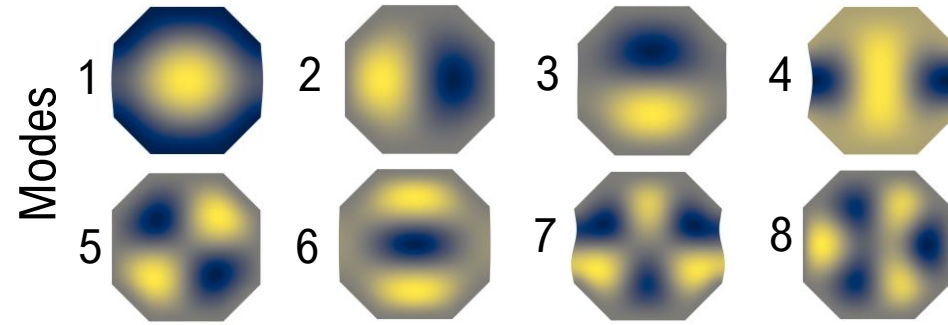
$$u_{k_y}(x, y) = e^{ik_y y} \phi_{\mathbf{k}}(x, y)$$



1<sup>st</sup> Brillouin zone

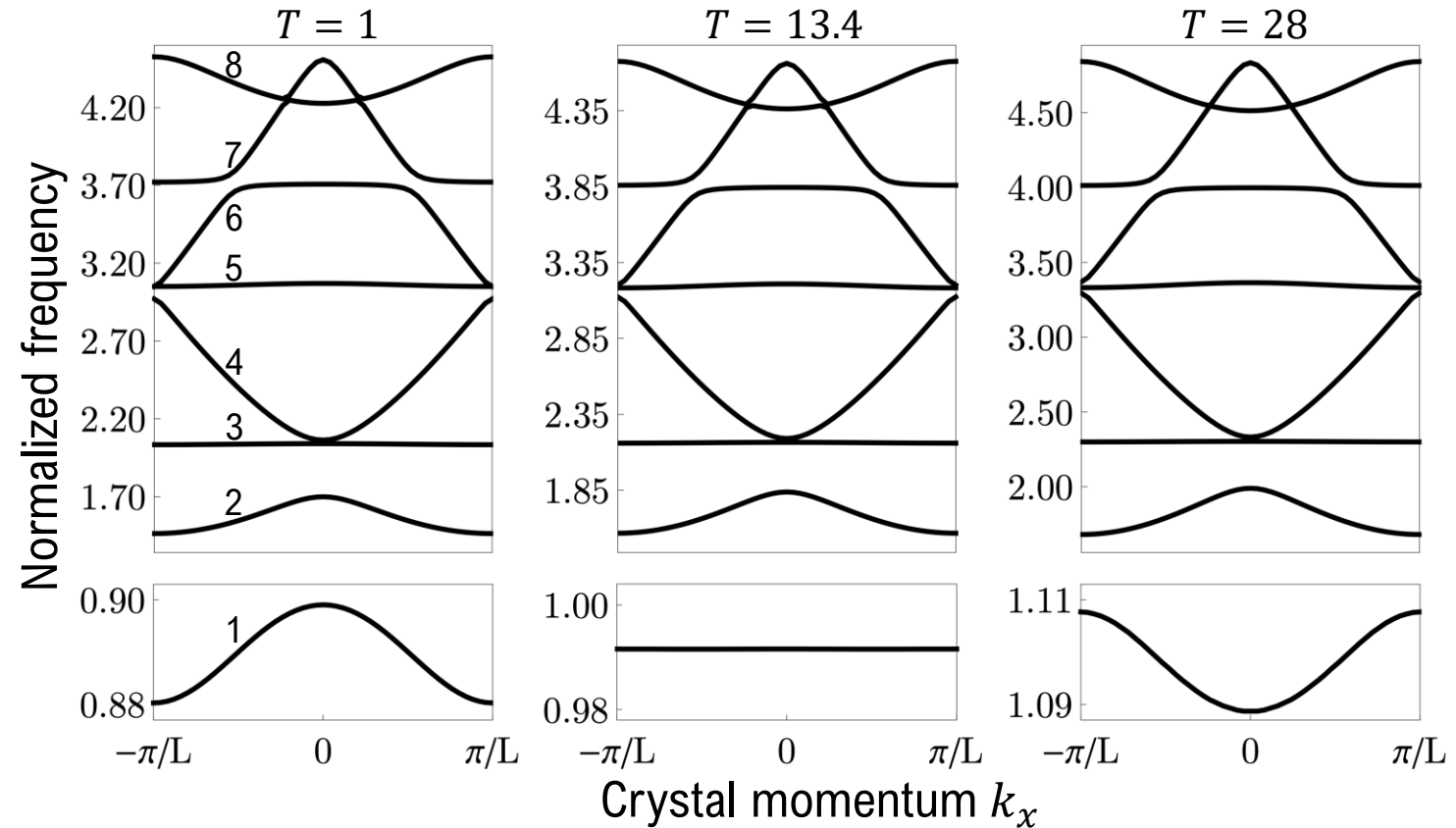
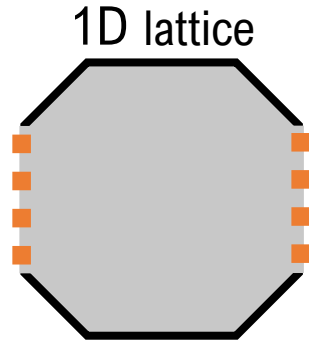


# Band structure calculation – 1D



Bloch periodicity:

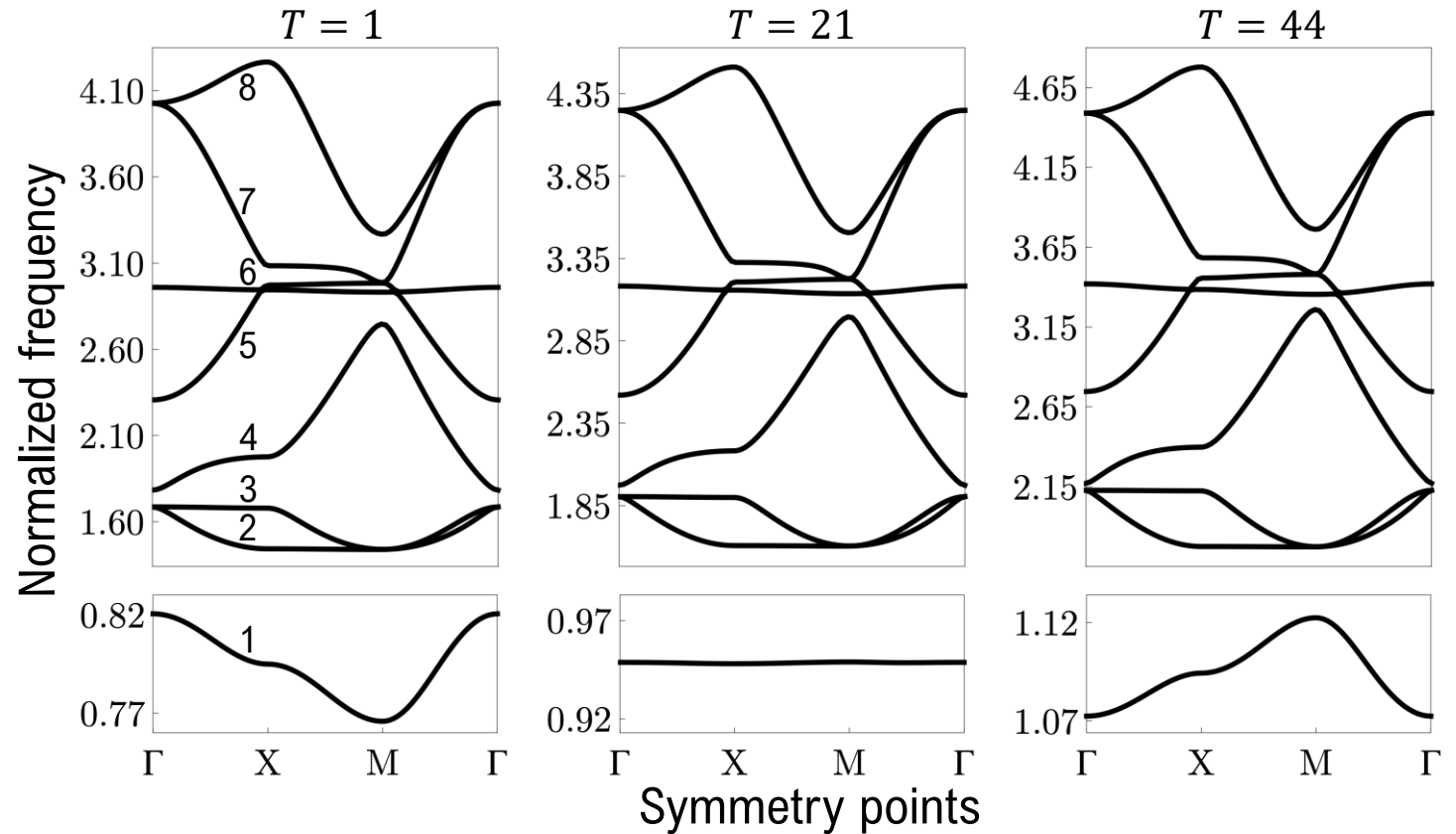
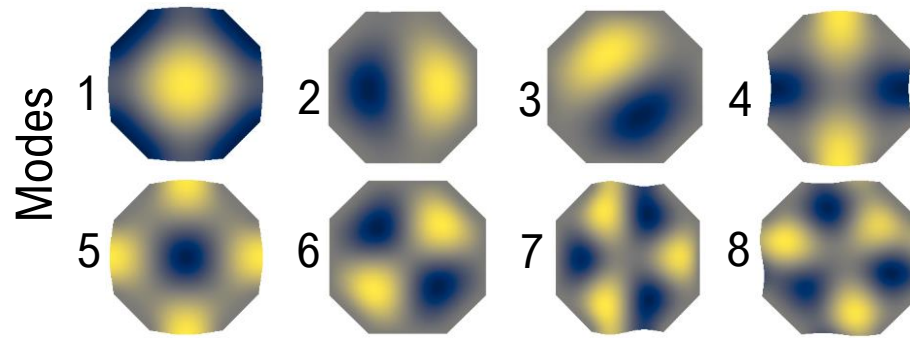
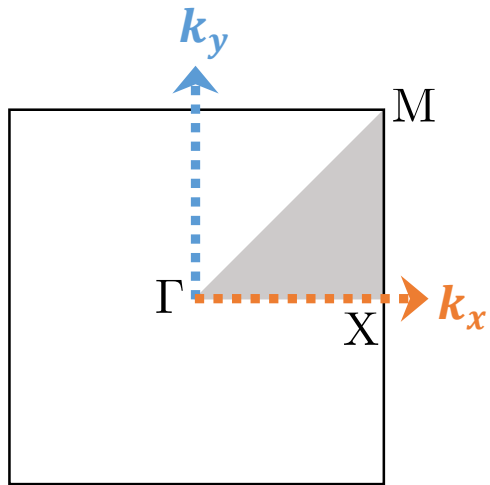
$$u_{k_x}(x, y) = e^{ik_x x} \phi_{k_x}(x, y)$$





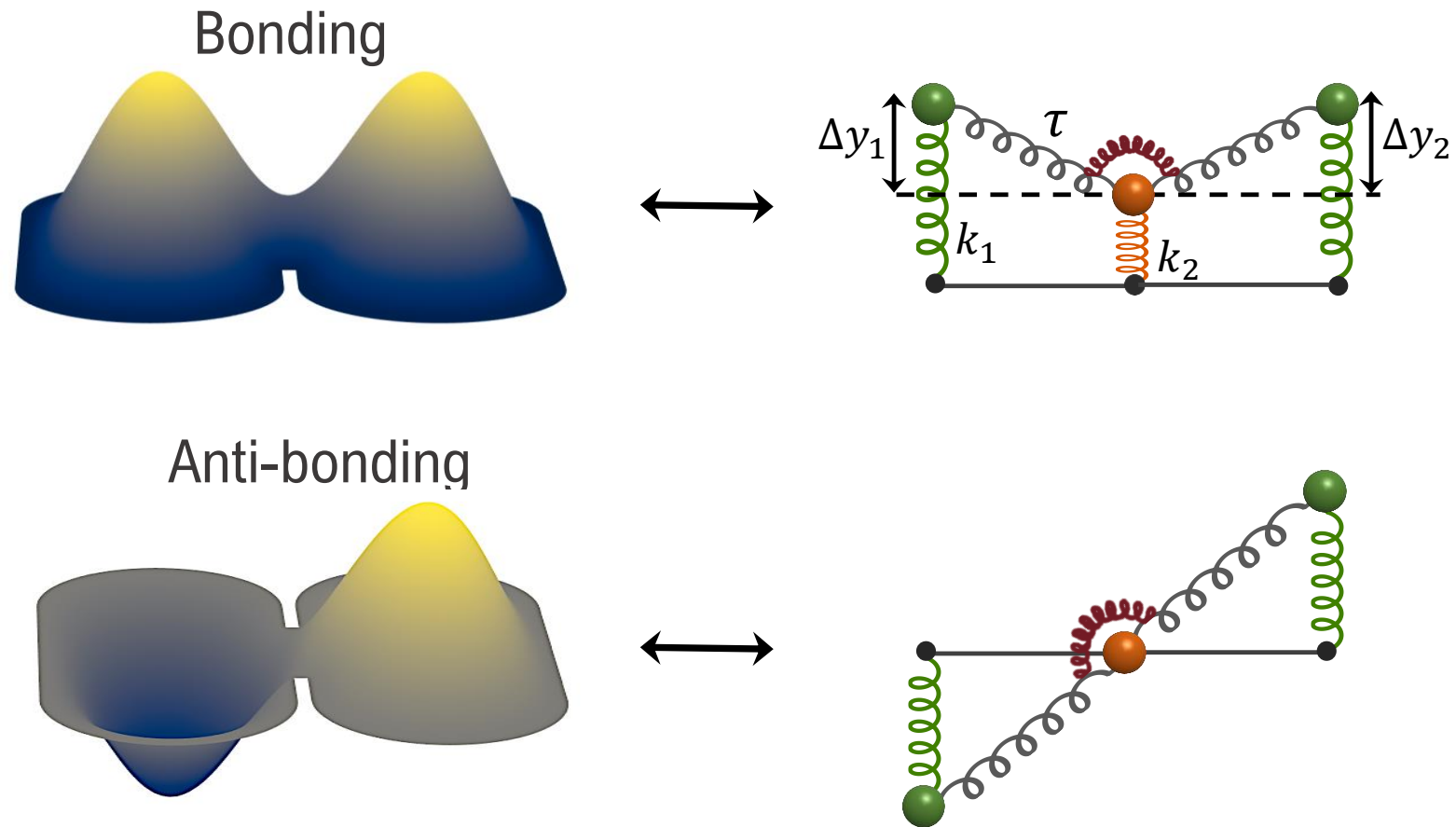
# Band structure calculation – 2D

1<sup>st</sup> Brillouin zone



# Mapping to a discrete spring mass model

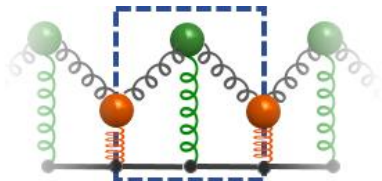
P. Karki and J. Paulose, *Physical Review Applied*, 2021.



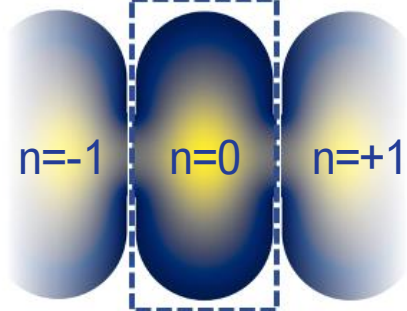
# Band structure of 1D chain

Periodic boundary condition

$$\mathbf{u}(n, t) = \mathbf{u}_0 e^{i(qn - \omega t)}$$



Unit cell

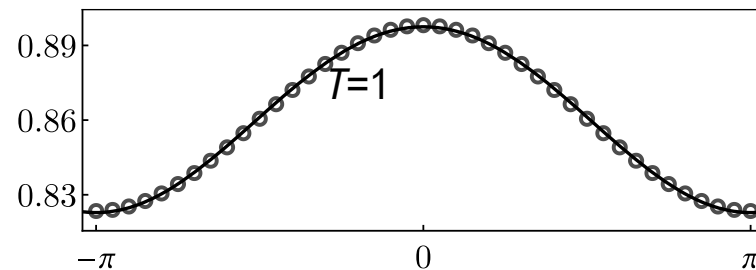


$$u_q(x, y) = \phi_q(x, y) e^{iqn}$$

Normalized frequency  $\omega/\omega_0$

○ COMSOL

— Theory

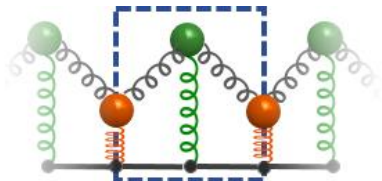


Quasi-momentum  $q$

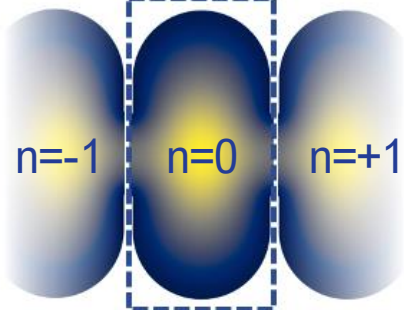
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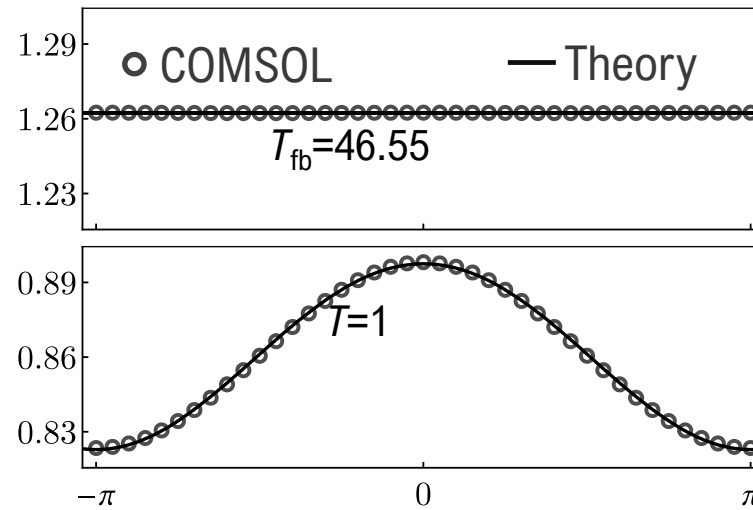


Unit cell



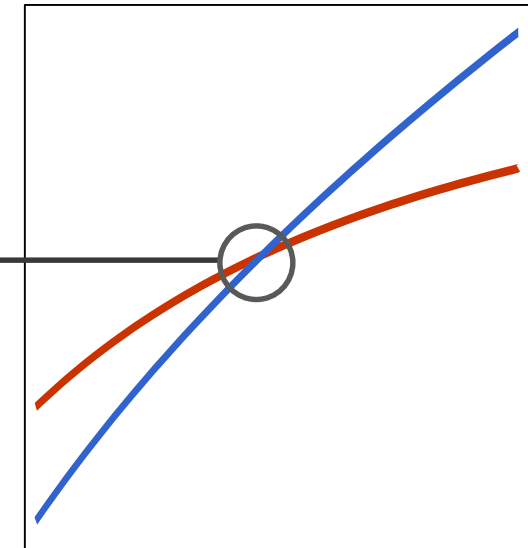
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Quasi-momentum  $q$

Mode-crossing  
 $\omega$  vs  $T$

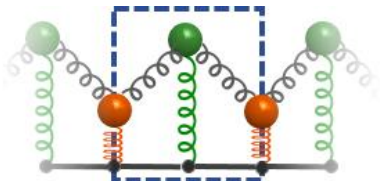


Tension  $T$

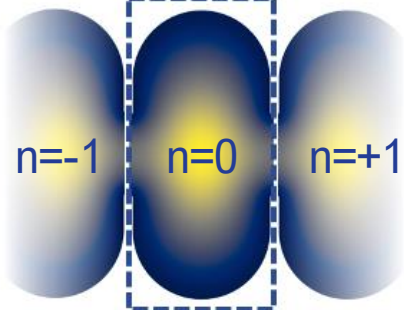
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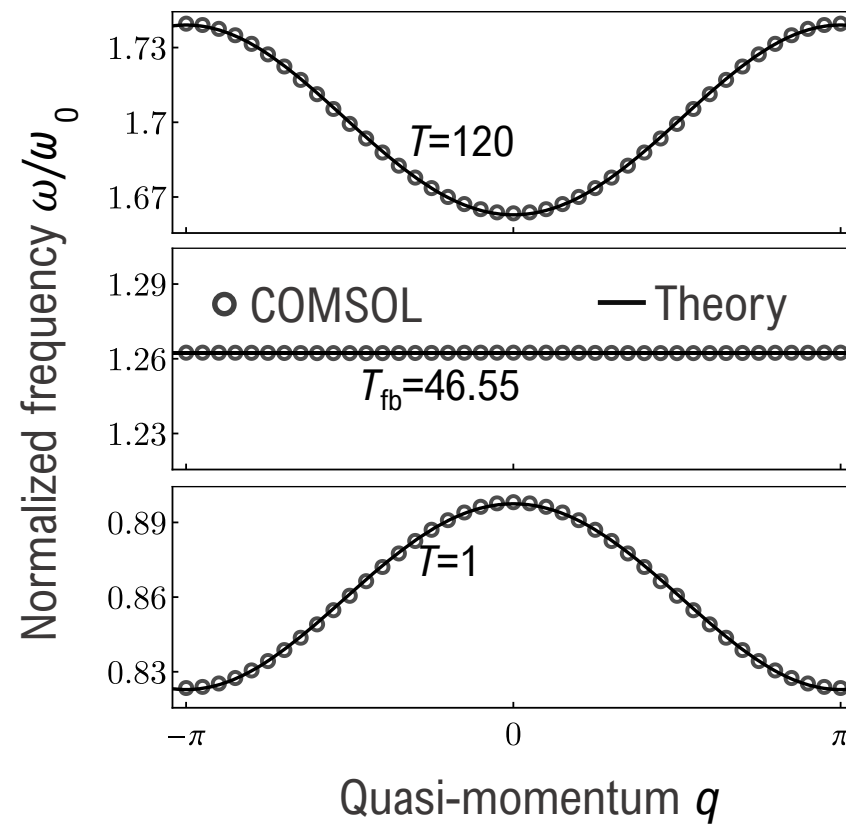
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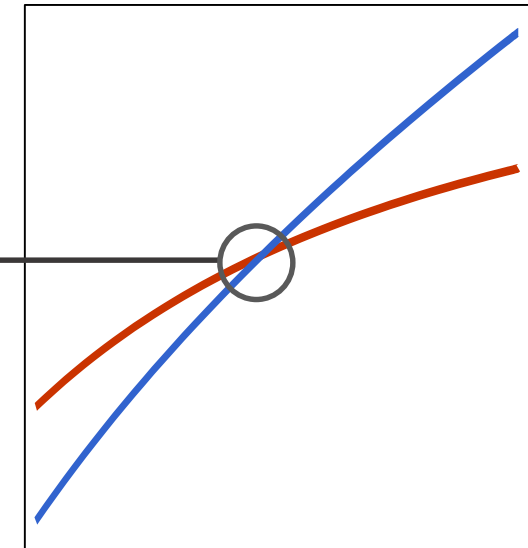
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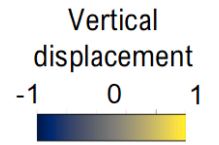
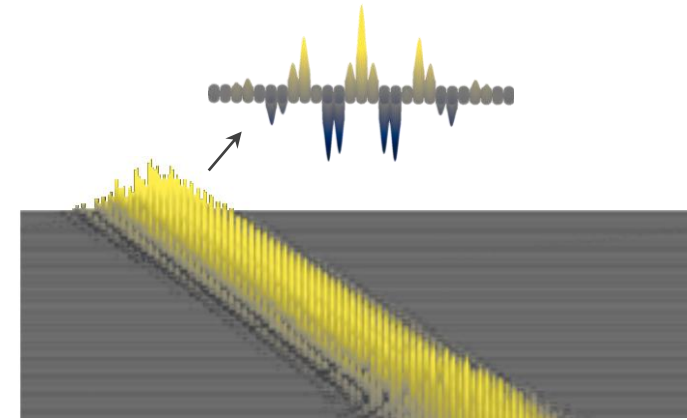
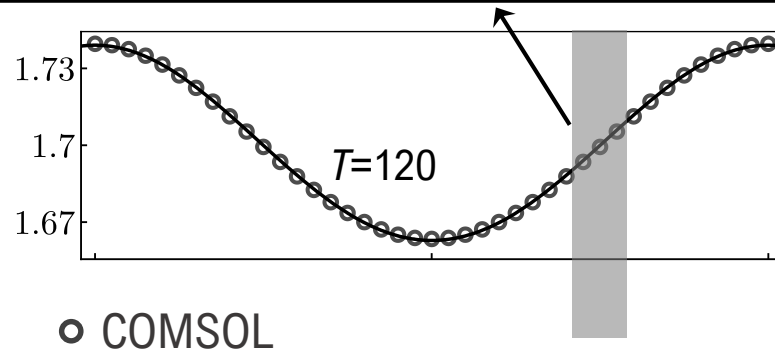
Mode-crossing  
 $\omega$  vs  $T$



Tension  $T$

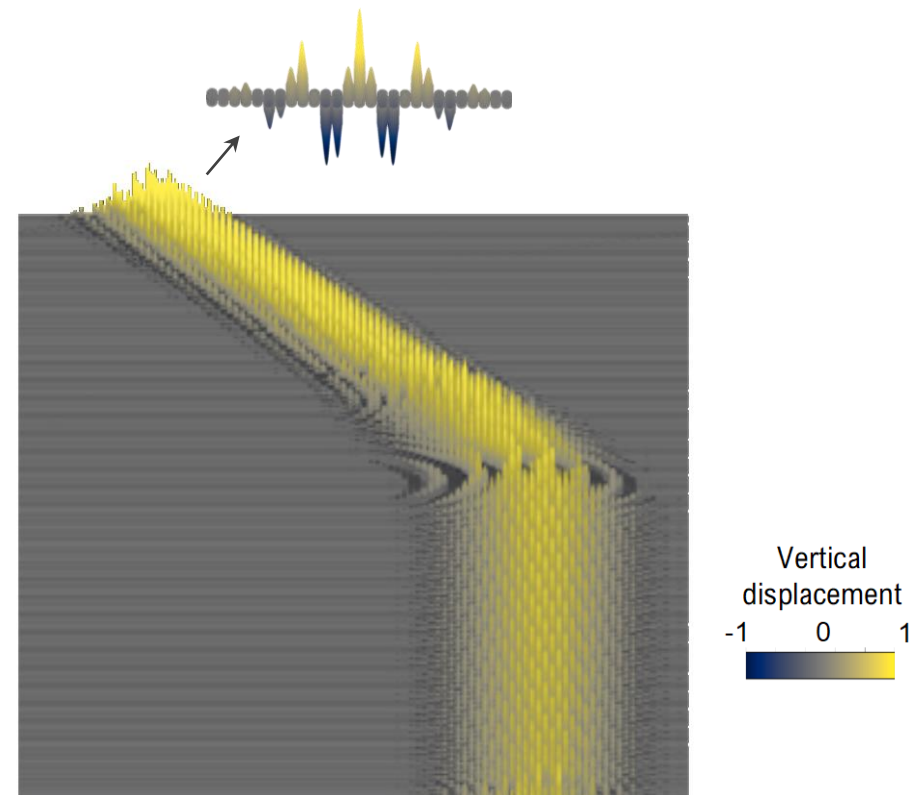
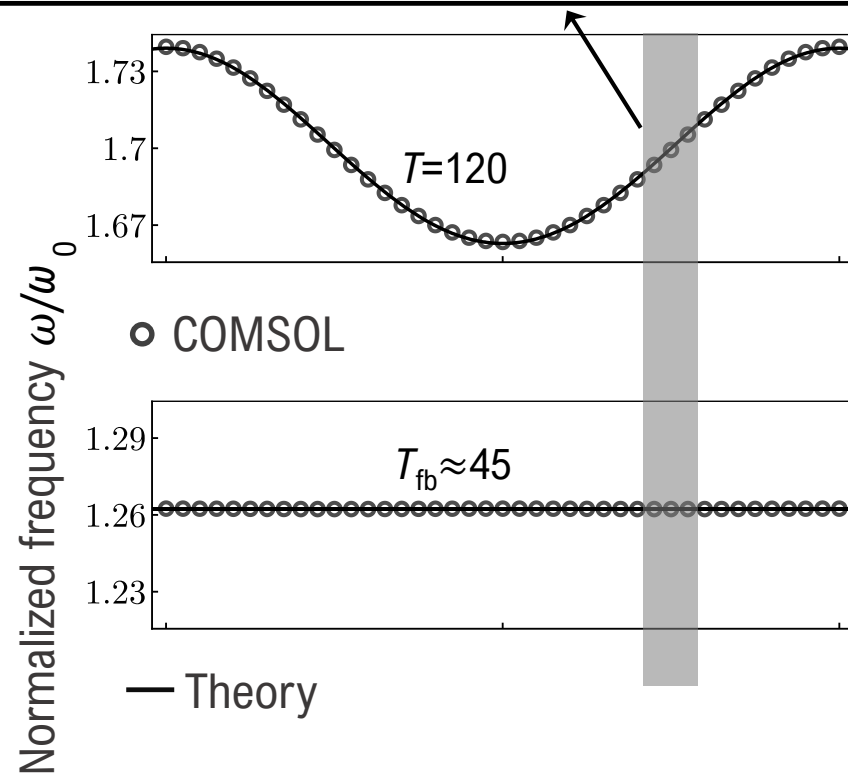
# Stopping and reversing sound

Gaussian wave-packet  $u(n, t) = \sum \phi_{q_x} e^{i(q_x n - \omega(q_x)t)} e^{-\left(\frac{q_x - q_0}{\Delta q_x}\right)^2}$



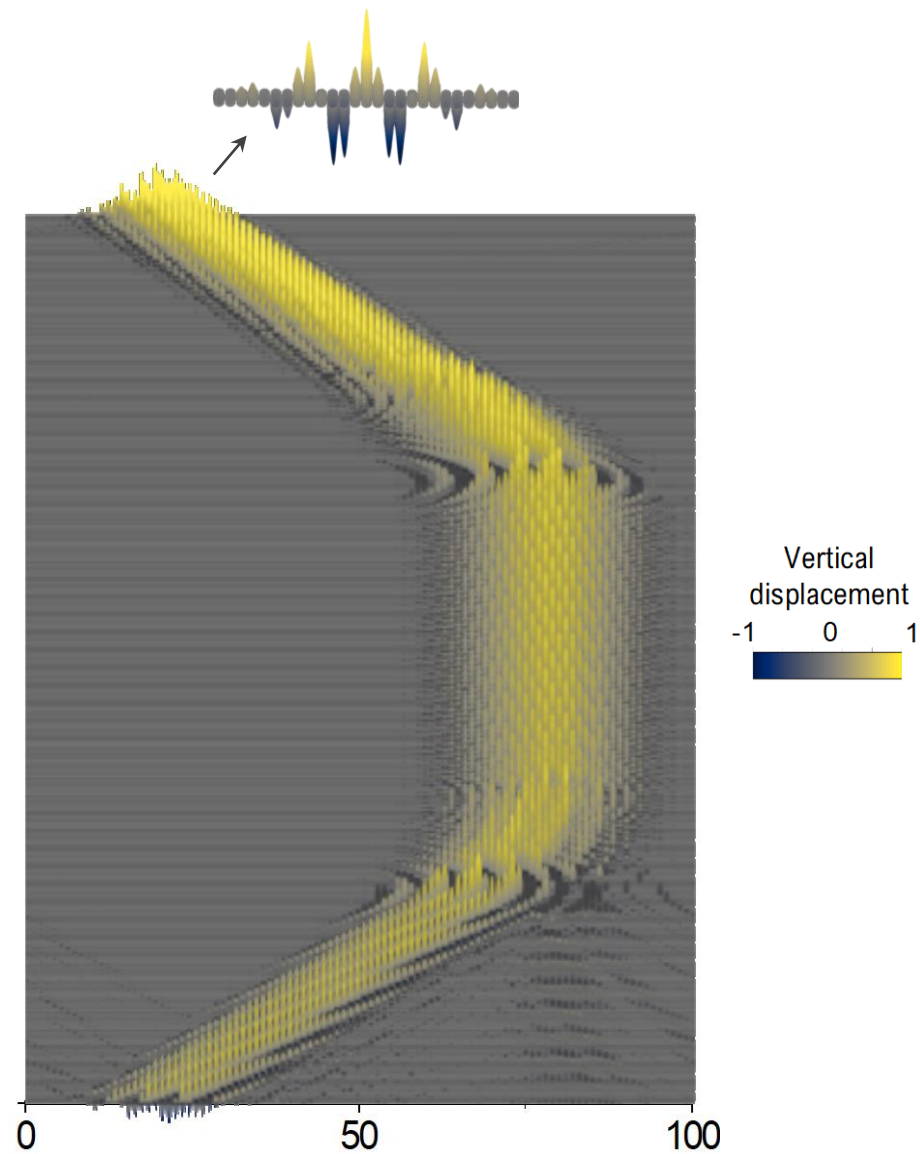
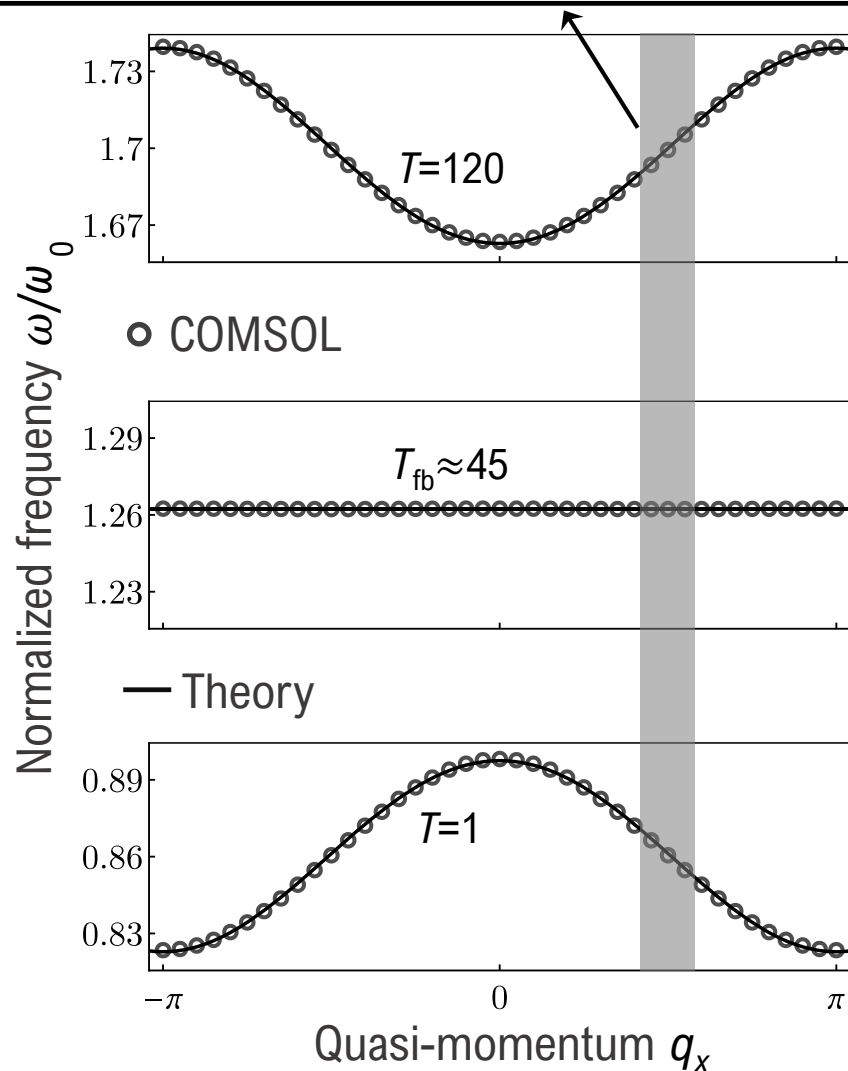
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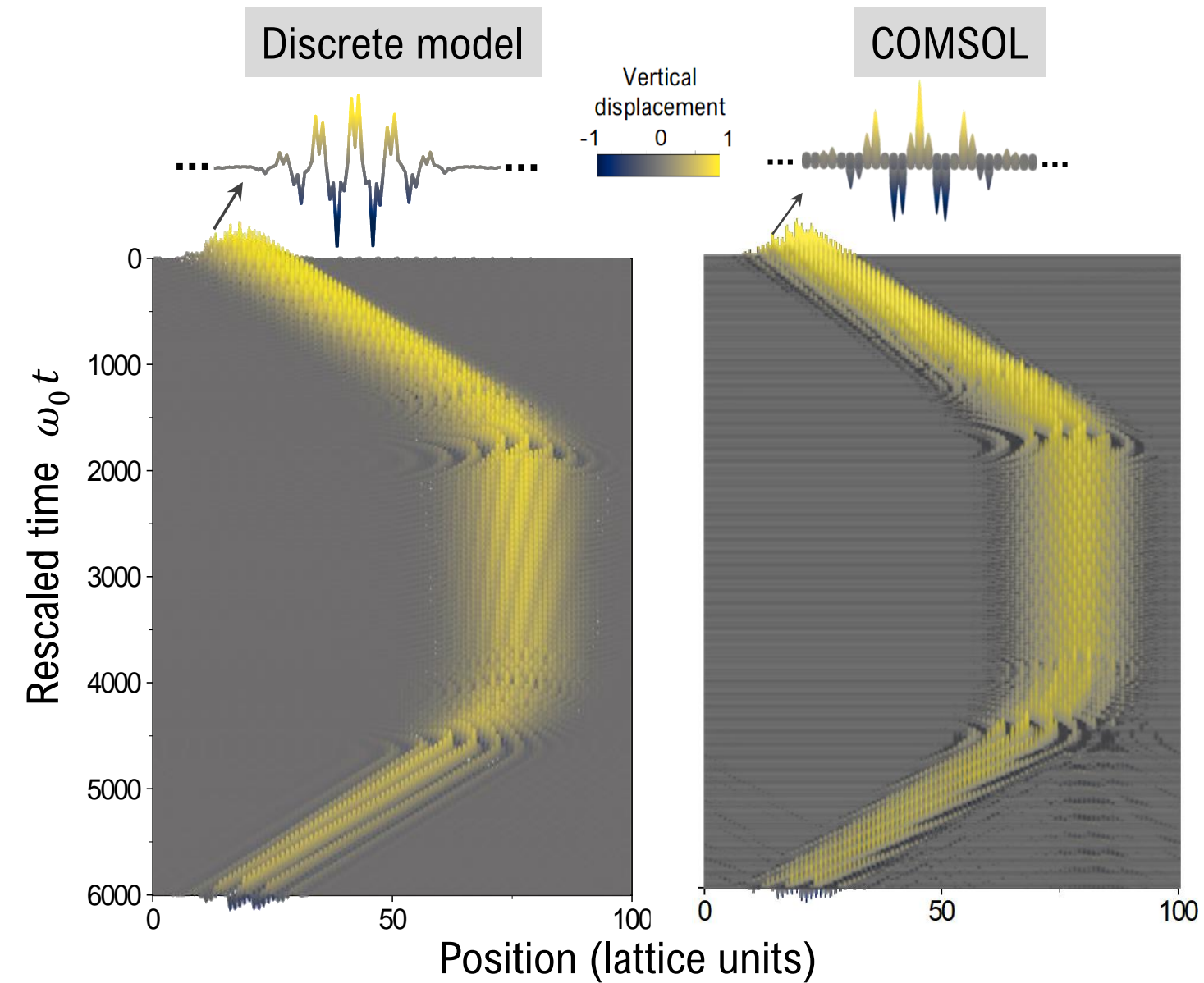




# Stopping and reversing sound



# Stopping and reversing sound



# Stopping and reversing sound

Discrete model

COMSOL

Vertical displacement  
-1 0 1

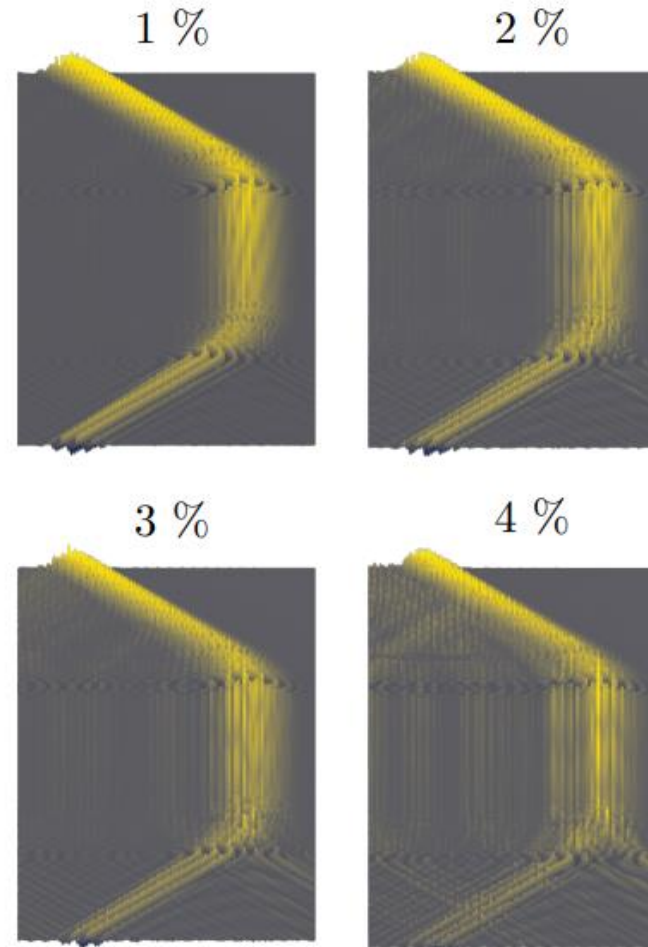
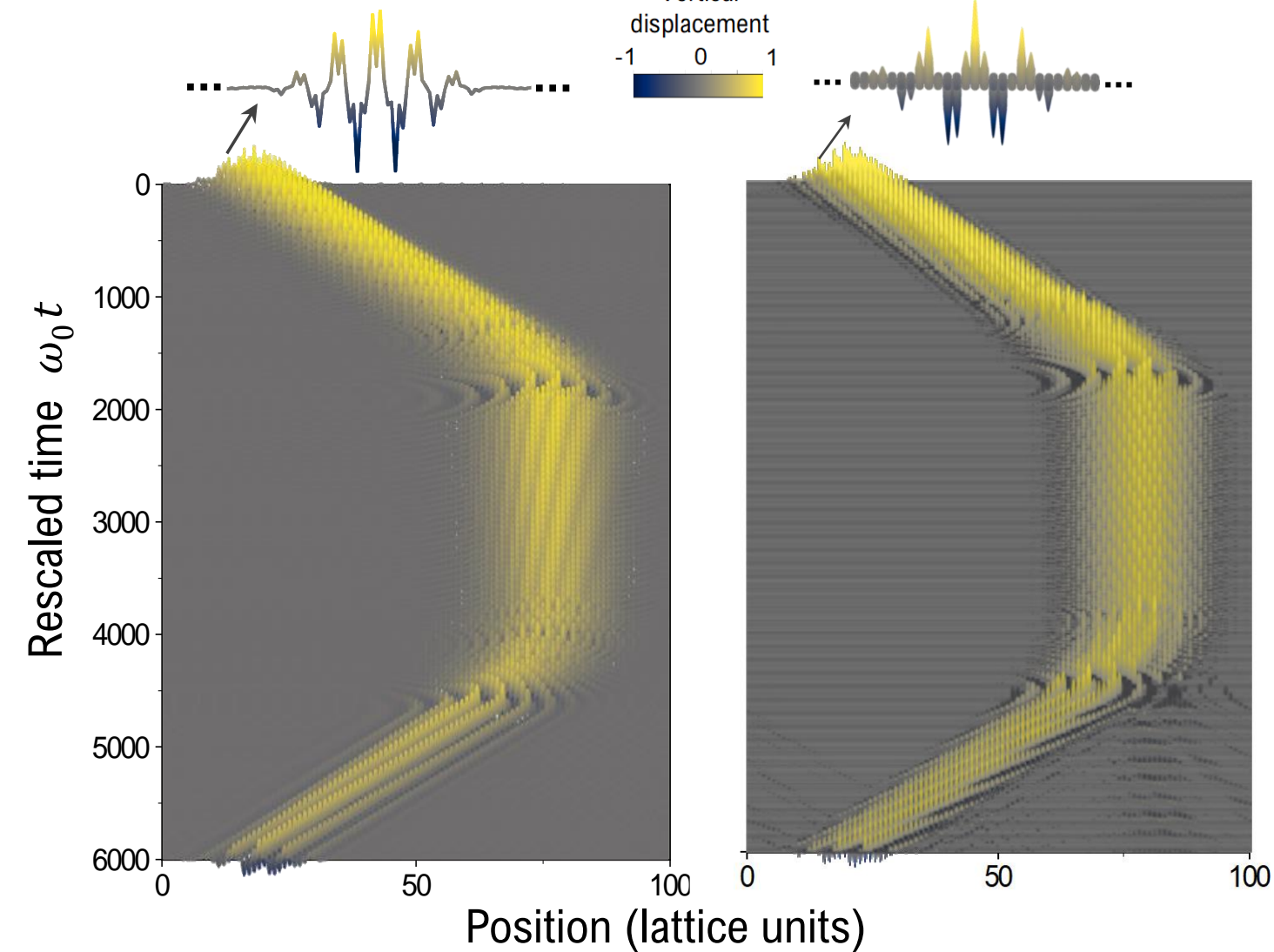
Disorder, % of  $\tau$  and  $k_1$ ,  
effects via discrete model

1 %

2 %

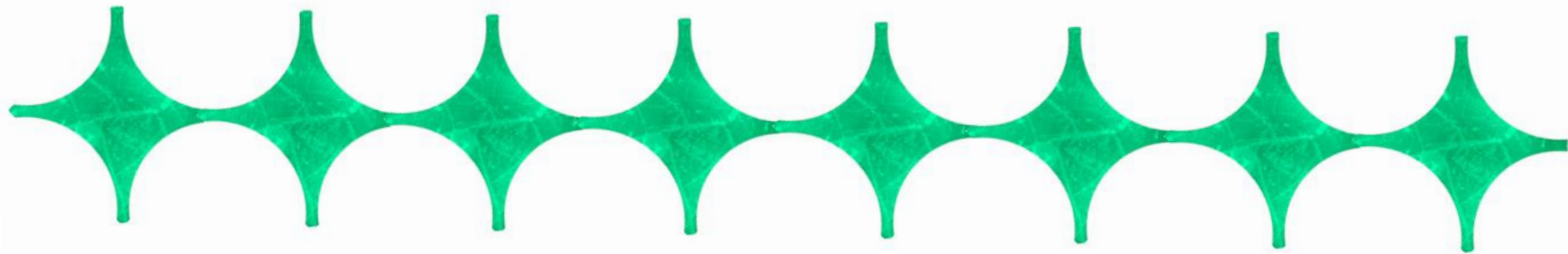
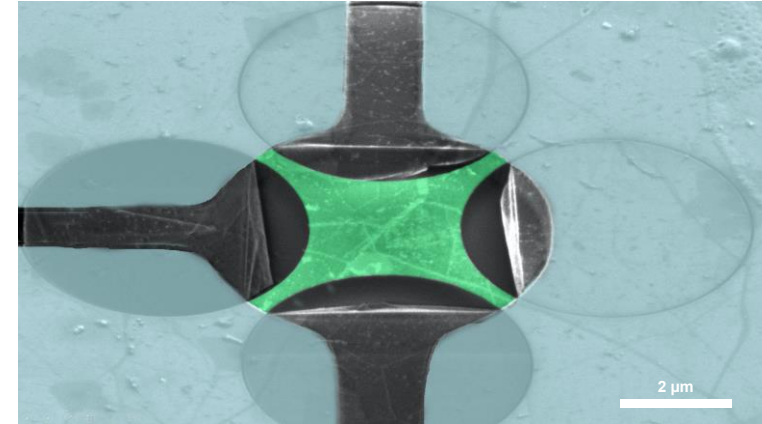
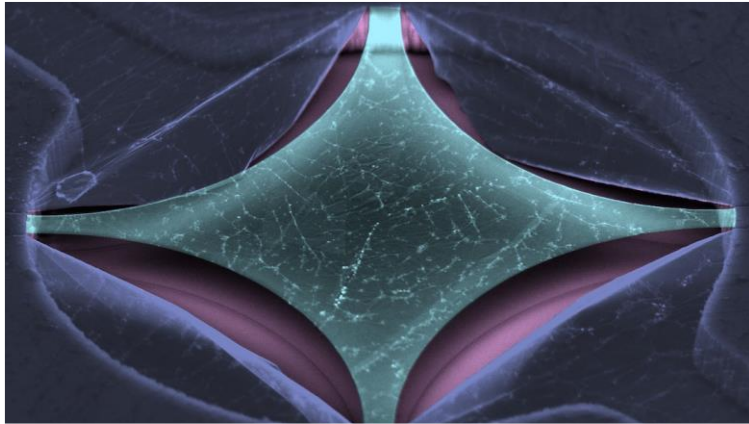
3 %

4 %



# Potential experimental system

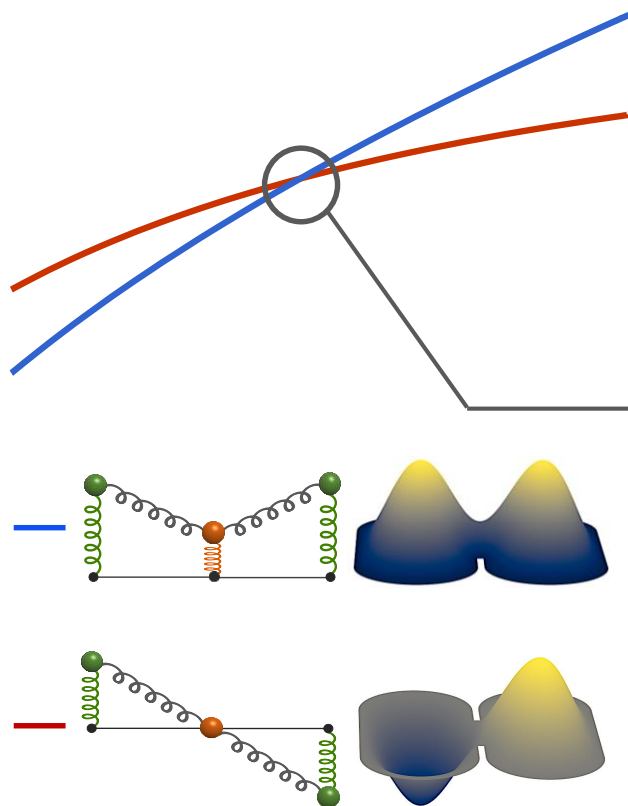
## Graphene



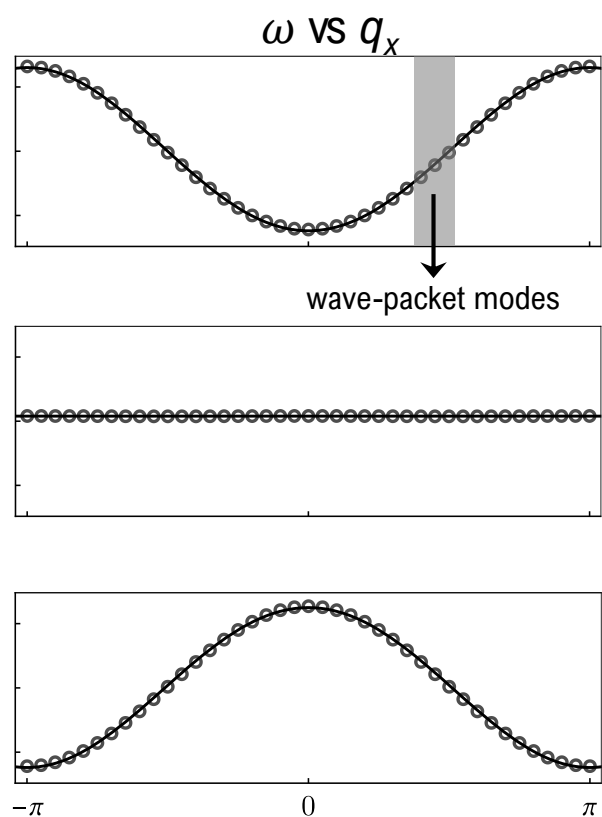
Benjamín Alemán group, University of Oregon

# Summary in 1D

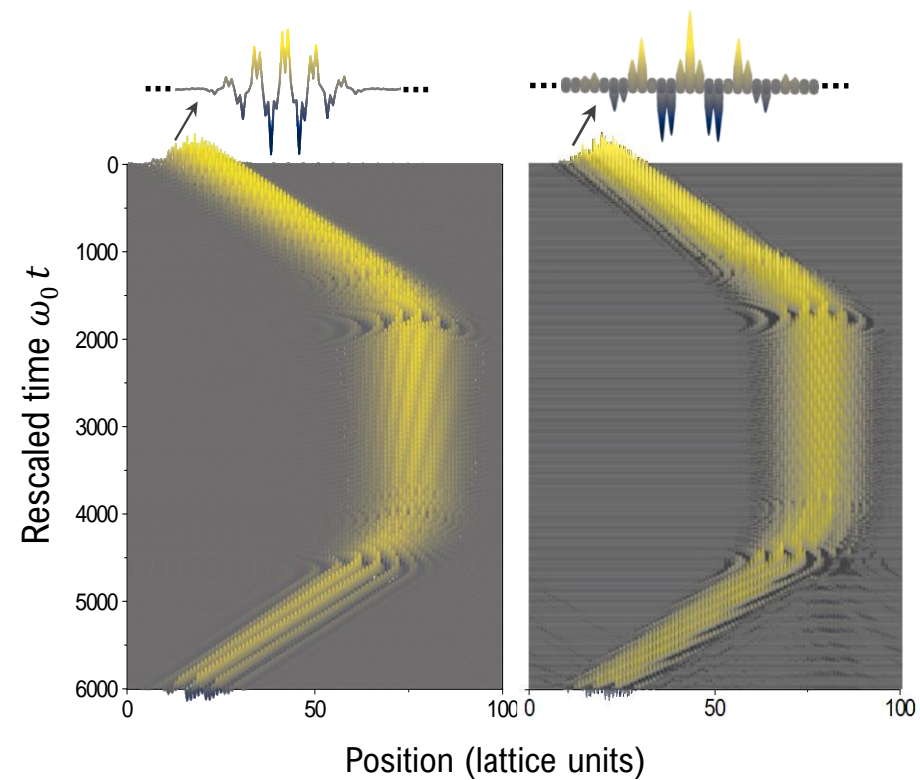
Vibrational modes crossing



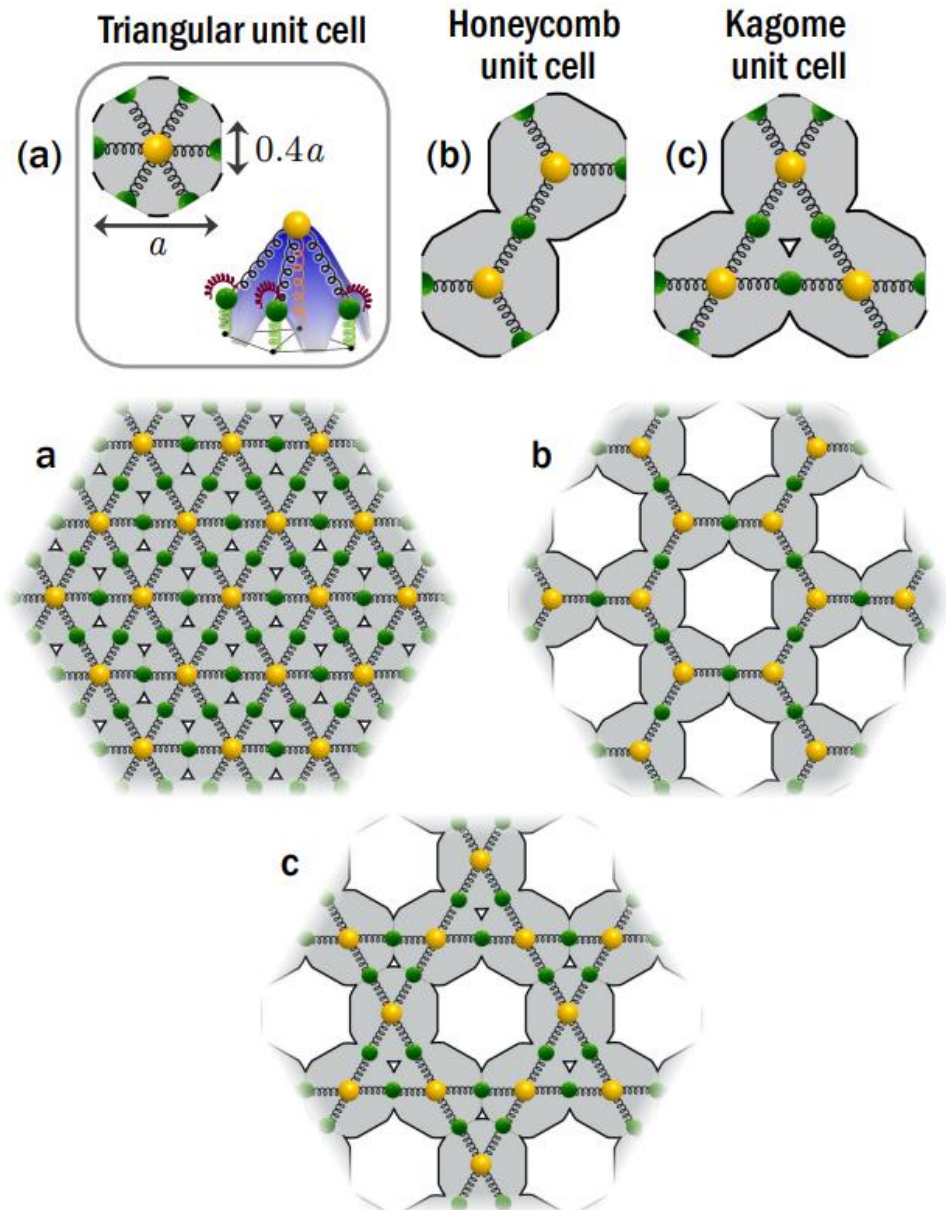
Changing dispersion character



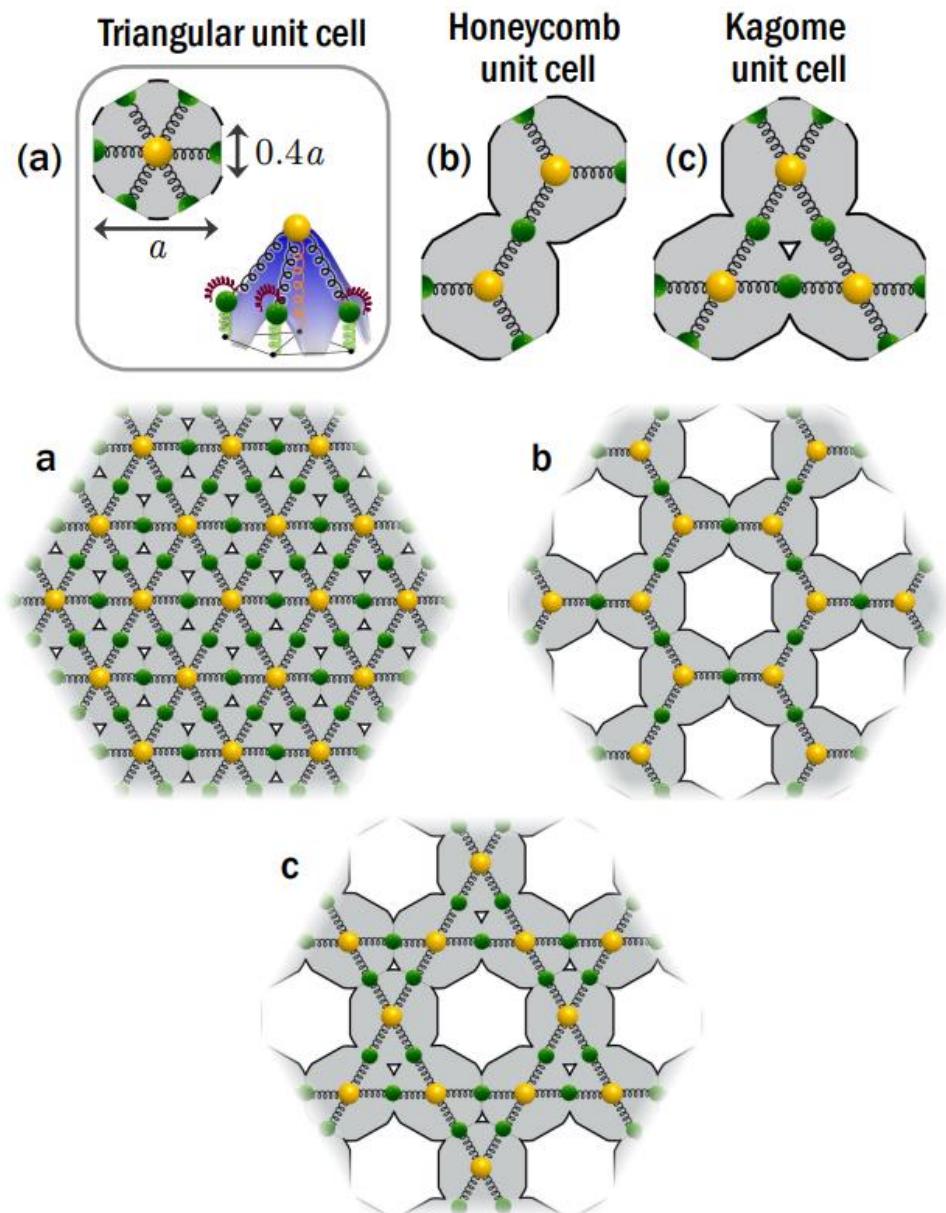
Stopping and reversing sound



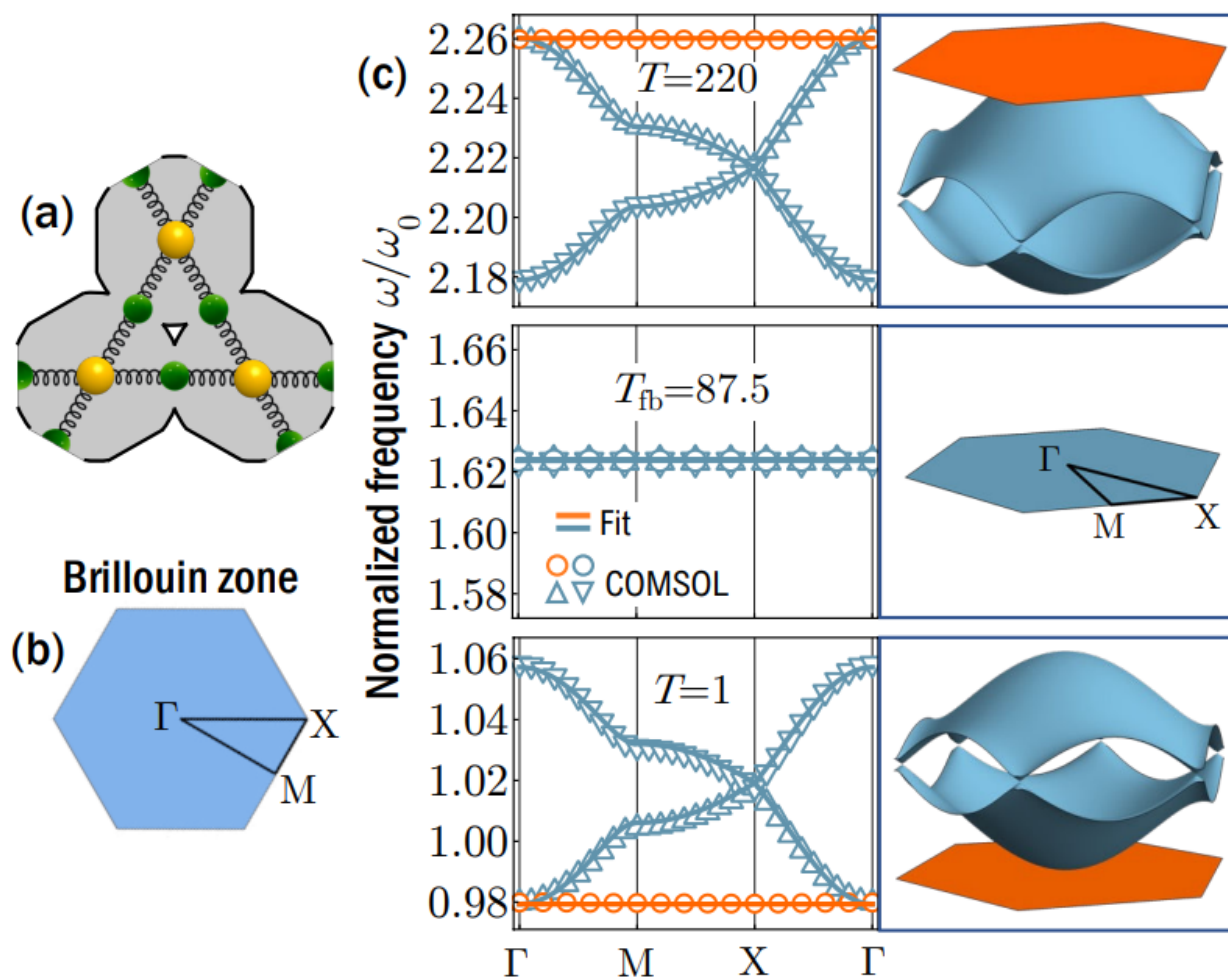
# Singular flat band in 2D



# Singular flat band in 2D



P. Karki and J. Paulose, *Physical Review Research*, 2023



# Conclusions

- We created a custom model for coupled thin-plate elastic resonators
- We modified the bonding character of the ground state of a fourth order coupled thin-plate elastic resonator system via prestress modulation
- We imported solutions from eigenvalue problem to time-dependent studies to perform dynamical simulations